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TITLE: QUANTITATIVE MEASUREMENTS OF LIGHT SCATTERING IN DECOMPOSING

**MASTER**

<sup>3</sup>He-<sup>4</sup>He LIQUID MIXTURES

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QUANTITATIVE MEASUREMENTS OF LIGHT SCATTERING IN DECOMPOSING  
 ${}^3\text{He}$ - ${}^4\text{He}$  LIQUID MIXTURES\*

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Measurements are presented of the angular distribution of light scattering by decomposing  ${}^3\text{He}$ - ${}^4\text{He}$  mixtures. It is found that the structure factor has a scaling form  $S(\kappa, t) = \kappa_m^{-3} F(\kappa/\kappa_m)$  where  $\kappa_m^{-1}$  is related to the size of a growing domain and  $F(\kappa)$  is time independent during the early stages of phase separation.

INTRODUCTION

We have made quantitative measurements of the angular distribution of light scattering by  ${}^3\text{He}$ - ${}^4\text{He}$  mixtures pressure quenched into the miscibility gap near the tricritical point. An earlier study[1] in  ${}^3\text{He}$ - ${}^4\text{He}$  showed that laser light scattered by the decomposing mixture is in the form of a collapsing ring similar to that observed in organic binary mixtures[2,3]. Recent Monte Carlo studies of Marro, Lebowitz, and Kalos[4] (MLK) have shown that the (normalized) structure factor  $S(\kappa, t)$  exhibits a scaling behavior in the homogeneous form

$$S(\kappa, t) = K^{-3} G(\kappa) F(\kappa/\kappa_m) \quad (1)$$

where  $K^{-3}$  is a measure of the time dependent domain size,  $\kappa_m^{-1}$  is related to time, and  $F(\kappa/\kappa_m)$  is independent of time. Experiments by photobric and scatter (PS) and  ${}^3\text{He}$ - ${}^4\text{He}$  data of Marro et al. [1] are used to validate the scaling form[1]. In this paper we show that a similar scaling behavior is valid for  ${}^3\text{He}$ - ${}^4\text{He}$  liquid mixtures. We also discuss some of the characteristics, generation, and differences in the angular distribution of light scattering in  ${}^3\text{He}$ - ${}^4\text{He}$  and organic binary mixtures.

EXPERIMENTAL

The experiments were carried out in a  ${}^3\text{He}$  cryostat with optical beam and optical fiber quartz windows. The apparatus was designed to give the minimum cell length that would allow the formation of the sample space. In this way the decomposition of the sample was not affected by the pressure and decompression process. A quartz fiber bundle with a core diameter of  $10^{-3}$  cm was used to monitor the rapid decompression of the mixture. Two thermometers located near the top and the bottom (inside the sample space) monitored the initial temperature and the slight temperature rise that accompanied decompression and phase separation.

The angular distribution of the scattered laser light ( $\lambda = 6328 \text{ \AA}$ ) was determined using a

Reticon photodiode array of 512 elements, the signals from which were stored in a transient recorder. Due to the large number of elements in the photodiode array, a continuous distribution of intensity could be easily obtained. Because each of the photodiode elements is rectangular (width = 13  $\mu\text{m}$ , height = 2.5 mm), the array was masked off to expose a wedge-shaped sensing area, thus preserving good resolution at low scattering angles. The output signals were subsequently plotted, smoothed, and appropriately corrected to yield the angular distribution of scattered light intensity  $I(\theta, t)$ . We have not as yet measured  $I(0, t)$  and thus cannot correct the data for effects of multiple scattering. However, because of the small scattering intensities, we expect such corrections to be small, even for our relatively long wavelength ( $6328 \text{ \AA}$ ).

RESULTS AND DISCUSSION

We now discuss the light scattering measurements made during the early stages of phase separation of an initially homogeneous mixture. Light scattering at a scattering vector  $\kappa = 0.910 \text{ \AA}^{-1}$  and at a quench pressure of  $3 \mu\text{atm}$  of forward light scattering was observed until the pressure in the cell dropped to one atmosphere. Then when a bright and sharply defined ring of light suddenly appeared (as is typically observed on a screen), this is consistent with the fact that light scattering is expected to be a non-linear effect the optical boundary is crossed. After its sudden appearance, the ring structure collapsed to the central beam which was approximately  $3 \text{ m}$ , which is about  $1000 \text{ \AA}$  in diameter. The initial distribution of the light was scattered light recorded by the photodiode array and the data was then quickly processed to give an intensity versus  $\theta$  plot of the ring which collapsed to the central laser beam. It is interesting to point out that the time evolution of the angular distribution of scattered light in  ${}^3\text{He}$ - ${}^4\text{He}$  has a striking resemblance to results from a Monte Carlo study of tricritical optical decomposition [5, 6].

metamagnet[5], where, for the case of the conserved order parameter, it is seen that the intensity distributions are more sharply peaked than those observed in organic binary mixtures.

In the cases of I-W and L-W, the maximum intensity  $I_m$  in the scattered light increased by almost three orders of magnitude before the ring finally collapsed to the central beam. In  $^3\text{He}$ - $^4\text{He}$ , the corresponding increase in  $I_m$  is only about a factor of 30. The wave vector  $\kappa_m$  ( $\kappa = [4\pi/\lambda]\sin\theta/2$ , where  $\theta$  is the scattering angle) corresponding to  $I_m$  was found to decrease in a similar manner to that observed in I-W and L-W, i.e.,  $\kappa_m \propto t^{-\phi}$ , where  $\phi$  gradually changes from  $3/2$  at early times to  $1$  at later times.

In Fig. 1 we have plotted the results of our measurements in terms of  $\kappa_m^3 I(\kappa, t)$  as a function of  $\kappa/\kappa_m$  for various times  $t$ . Whereas MIF found that Eq. (1) was best satisfied when  $K$  was taken as the first moment,  $k_1$ , of the structure factor, we follow Ref. 7 and identify this parameter with  $\kappa_m$ . Note the sharpness of the peaks in the scaled curves (Fig. 1). The overall shape of the curves is quite similar to that presented by MIF (see Fig. 2 of Ref. 4). It is obvious from Fig. 1 that the function  $I(\kappa/\kappa_m)$  is reasonably independent of time, at least for values of  $\kappa/\kappa_m \geq 1$ . To the left of the peak in Fig. 1, the data show a slight trend with time, suggesting that at the earlier times  $\kappa_m$  does not have a well-defined value. This would be consistent with the findings of MIF, who claims that the appearance of a first-order scattering ring is a consequence of the scattering beam scattering off a region of finite size. The appearance of the peak at  $\kappa/\kappa_m = 1$  is a small background of intensity at small scattering angles. The correction affects the data for the earliest times somewhat more than at other times. In addition, a slight expansion of the scattering angle would affect the data at large scattering angles more than at small angles. The authors of Ref. 4 claim that the curves in Fig. 1 do not satisfy Eq. (1) since the curves are not symmetric about  $\kappa/\kappa_m = 1$ . However, the present data are not essential to test Eq. (1) for this reason.

In order to compare our results with those of Ref. 4, we have expressed our data in terms of the normalized function  $I(\kappa/\kappa_m)$ . Curve 1 in Fig. 1 is the first-order scattering ring at a scattering length  $\lambda$  of  $1.0 \mu\text{m}$ . The data were plotted for the purpose of testing Eq. (1) and were not the subject of the present study. The view shown in Fig. 1 is that they, with the possible exception of the first time frame at  $t = 0.5$   $\mu\text{sec}$ , shown in Fig. 1, our data confirm the findings of Ref. 4. That the scaling law, Eq. (1), is valid for early times  $t$ , even over a wide range of  $\lambda$ ,

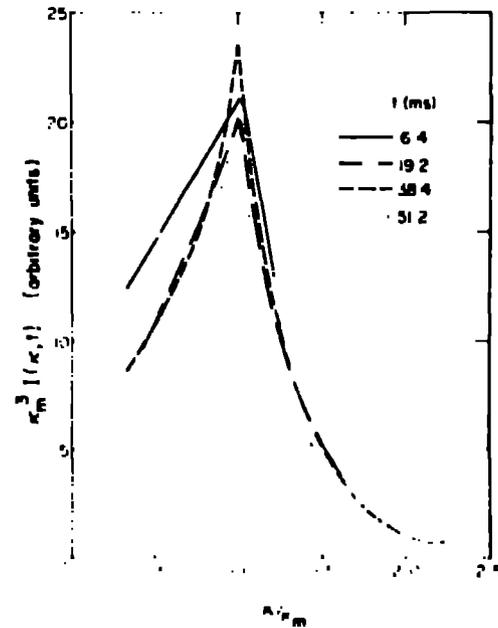


Fig. 1. Smoothed intensities of scattered laser light, scaled by  $\kappa_m^3$ , vs  $\kappa/\kappa_m$  for a mixture of liquid  $^3\text{He}$ - $^4\text{He}$  at composition  $X = 0.673$  and initial temperature  $0.816$  K, following a pressure quench from  $0.4$  atm to  $0.2$  atm. With the exception of the earliest time frame, the data appear to scale with the asymptotic expression  $I(\kappa/\kappa_m) \propto \exp(-\kappa/\kappa_m)$  for  $\kappa/\kappa_m \geq 1$ .

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